



## Real Spectra in Logarithmic model PT-symmetry operators: Iso-spectra in Logarithmic PT-symmetry

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**Abstract.** We reflect real spectra of new logarithmic model PT-symmetry operators with singular and non-singular in nature. We also notice the iso-spectral nature between inverted and non-inverted PT-symmetry potentials. Present numerical result give good agreement with previous results.

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### 1. Introduction

Real spectra in quantum operators are confined to Hermiticity ( $H = H^\dagger$ ) as well as PT-symmetry [1] ( $[H, PT] = 0$ ). Here  $P$  stands for parity operator having the properties:  $PxP^{-1} = -x$ ;  $PpP^{-1} = -p$ . Similarly  $T$  stands for the time reversal operator having the properties  $TxT^{-1} = x$ ;  $TpT^{-1} = -p$  and  $TT^{-1} = -i$ . In the Hermiticity (more precisely self-adjoint operator), it is possible to find two Hamiltonians, which are iso-spectral to each other [2]. For example [2,3]

$$h^{(1)} = p^2 + V_0(1 - e^{2|x|/a}) \quad (1)$$

and

$$h^{(2)} = p^2 + v_0(1 - e^{-2|x|/a}) \quad (2)$$

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If one clearly analyzes one is scattering in nature and the other is confining nature [2,3].

Till now no such models are reflected in nature. In a recent paper Bender et al [4] have suggested a new class of logarithmic PT-symmetry potentials as

$$h_1 = p^2 + x^4 \log(ix) \quad (3)$$

$$h_2 = p^2 - x^4 \log(ix) \quad (4)$$

$$h_3 = p^2 - x^4 \log(x^2) \quad (5)$$

and reflected energy spectrum of  $H_1$  only. Further authors reported analytical calculation of energy level using WKB approach[4]

$$\frac{E_n}{[\log E_n]^{1/3}} \sim \left[ \frac{\Gamma(7/4)(n + 1/2)\sqrt{\pi}}{\Gamma(5/4)\sqrt{2}} \right]^{4/3} \quad (6)$$

does not give encouraging results when compared with numerical results. This motivates the present author to calculate spectra of  $H_{1,2}$  and suggest new models on logarithmic potentials. Apart from this aim is to find out whether iso-spectral Hamiltonians are possible in PT-symmetry operators.

## 2. Logarithmic new models

Here we consider different models as follows

**Quadratic Logarithmic model**  $V(x) = -x^2 \log(\frac{i}{x})$

The Hamiltonian considered here is

$$H_1^{quadratic} = p^2 - x^2 \log(\frac{i}{x}) \quad (7)$$

**Quartic inverted model**  $V(x) = -x^4 \log(\frac{i}{x})$

Here we consider the Hamiltonian

$$H_2^{quartic} = p^2 - x^4 \log(\frac{i}{x}) \quad (8)$$

Below we present few energy levels

## 3. Logarithmic models

Here we consider the recently proposed models

$$h_1 = p^2 + \lambda x^4 \log(ix) \quad (9)$$

$$h_2 = p^2 - \lambda x^4 \log(ix) \quad (10)$$

$$h_3 = p^2 - \lambda x^4 \log(x^2) \quad (11)$$

and present complete spectra in table-2 for  $\lambda = 1$ .

TABLE.I: PT-symmetry inverted logarithmic model potentials

n	$H_1$ Present	$H_2$
0	1.326 591 6	1.249 087 3
3	9.173 294 3	13.738 280 4
6	17.734 002 2	31.665 810 8
9	26.633 530 1	52.993 926 2
12	76.113 329 2	76.974 762 5

#### 4. Method of calculation

Here we use matrix diagonalisation method [5] to solve the eigenvalue relation

$$H|\Psi\rangle = E|\Psi\rangle \quad (12)$$

where

$$|\Psi\rangle = \sum A_m |m\rangle \quad (13)$$

Here  $|m\rangle$  satisfy the relation

$$[p^2 + x^2]|m\rangle = (2m + 1)|m\rangle \quad (14)$$

Numerical results obtained using MDM are tabulated in table.1.

#### 5. Conclusion

In this report, we present numerical convergent energy levels of new model PT-invariant Hamiltonians using matrix diagonalisation method[5]. Further we feel the present method can be used confidently to realize real spectra study in similar Hamiltonians of interest. Lastly we the spectra of  $H_2$  and  $h_1$  are the same. In brief

$$V(x) = x^4 \log(ix) \rightarrow V(x) = -x^4 \log\left(\frac{i}{x}\right) \quad (15)$$

Hence these two potentials can be considered as iso-spectral models in PT-symmetry. Lastly we do not find any numerical results to present in table-1 for a comparison with the present numerals. Interested readers can consider other values of  $\lambda$ .

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TABLE.2: PT-symmetry logarithmic model

n	$h_1$ Present	$h_1$ Previous [4]	Previous (WKB)[4]
0	1.249 08	1.249 09	0.546 27
3	13.738 27	13.738 3	7.314 80
6	31.665 82	31.665 8	16.697 9
9	52.993 79	52.993 9	27.695 6
12	76.976 08	76.974 8	39.932 4
n	$h_2$ Present	Previous[4]	Previous(WKB)[4]
0	0.109 1		
1	6,959 6		
2	8.257 1		
3	18.039 4		
n	$h_3$ Present	Previous[4]	Previous(WKB)[4]
0	0.025 4		
1	4.977 7		
2	9.237 1		
3	16.478 6		

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