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# Boundary sentinel with given sensitivity in population dynamics problem and parameters identification 

Mifiamba Soma ${ }^{1}$, Somdouda Sawadogo ${ }^{2, *}$<br>${ }^{1}$ Département de Mathématiques et Informatique, Université Joseph Ki-ZERBO, Ouagadougou, Burkina Faso.<br>${ }^{2}$ Département de Mathématiques, Institut des Sciences, Ouagadougou, Burkina Faso.


#### Abstract

The notion of sentinels with given sensitivity was introduced by J.L.Lions [11] in order to identify parameters in the problem of pollution ruled by a parabolic equation. He proves that the existence of such sentinels is reduced to the solution of exact controllability problem with constraints on the state. In population dynamics model, we reconsider this notion of sentinels in a more general framework. We prove the existence of the boundary sentinels by solving a boundary null-controllability problem with constraint on the control. Our results use Carleman inequality which is adapted to the constraint.


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Key Words and Phrases: Population dynamics, optimal control, controllability, sentinels, Carleman inequality

## 1. Introduction

The notion of sentinel was introduced by J.L.Lions to study systems with incomplete data [11]. The notion permits us to distinguish and to analyse two types of incomplete data: the so-called pollution terms at which we look for information, independently of the other type of incomplete data which is the missing terms and that we do not want to identify.

Typically, the Lions's sentinel is a functional defined on an open set $O$ where we consider three functions: the "observation" $y_{\text {obs }}$ corresponding to measurements, a given "mean" function $h_{0}$, and a control function $w$ to be determined.

Let us remind that Lions's sentinel theory [11] relies on the following three features: the state equation $y$ which is governed by a partial differential equation, the observation system and some particular evaluation function: the sentinel itself. More precisely, we consider a linear model (1) describing the dynamics of population with age dependence, spatial structure with incomplete data.

[^0]Let $\Omega$ be an open and bounded domain of $\mathbb{R}^{N}, N \in\{1,2,3\}$, with boundary $\Gamma$ of $C^{\infty}$. For the time $T>0$ and the life expectancy of an individual $A>0$, we set $U=(0, T) \times(0, A), Q=U \times \Omega, Q_{A}=(0, A) \times \Omega, Q_{T}=(0, T) \times \Omega, \Sigma=U \times \Gamma$, $\Sigma_{1}=U \times \Gamma_{1}$, where $\Gamma_{1}$ is a non-empty open subset of $\Gamma$. Then consider the following two stroke problem:

$$
\left\{\begin{array}{cccc}
\frac{\partial y}{\partial t}+\frac{\partial y}{\partial a}-\Delta y+\mu y & = & 0 & \text { in }  \tag{1}\\
y(0, a, x) & = & y^{0}+\tau \hat{y}^{0} & \text { in } Q_{A} \\
y(t, 0, x) & = & \text { in } Q_{T}
\end{array}\right\} \begin{array}{ccc}
\int_{0}^{A} \beta(t, a, x) y(t, a, x) d a & & \\
y & =\left\{\begin{array}{ccc}
\xi+\sum_{i=1}^{M} \lambda_{i} \hat{\xi}_{i} & \text { on } & \Sigma_{1} \\
0 & \text { on } \Sigma \backslash \Sigma_{1}
\end{array}\right. &
\end{array}
$$

where :

- $y(t, a, x)$ is the distribution of a-year old individuals at time $t$ at the point $x \in \Omega$.
- $\beta(t, a, x) \geq 0$ and $\mu(t, a, x) \geq 0$ are respectively the natural fertility and the natural death rate of age a at time t and position $x \in \Omega$.
- Thus, the formula $\int_{0}^{A} \beta(t, a, x) y(t, a, x) d a$ denotes the distribution of newborn individuals at time $t$ and location $x$.
- The boundary condition is unknown on a part $\Sigma_{1}$ of the boundary and represents a pollution with a structure of the form $\xi+\sum_{i=1}^{M} \lambda_{i} \hat{\xi}_{i}$. In this structure, the functions $\xi$ and $\hat{\xi}_{i}, i=1, \ldots M$ are known whereas the real $\lambda_{i}, i=1, \ldots M$ are unknown.
- The initial distribution of individuals is unknown and its structure is of the form $y^{0}+\tau \hat{y}^{0}$ where the function $y^{0}$ is known and the term $\tau \hat{y}^{0}$ is unknown.

System (1) is a system with incomplete data because the information on the boundary condition as well as on the initial condition are partially or completely unknown. Here, the pollution is isolated on the boundary $\Gamma \backslash \Gamma_{1}$. The missing term is located in the initial conditions. In what follows, we assume as in [8] that:
$(H 1):\left\{\begin{array}{c}\beta \in L^{\infty}(Q), \quad \beta(t, a, x) \geq 0 \text { a.e. in } Q ; \\ \sup _{\substack{(t, x) \in] 0, T[\times \Omega}} \int_{j 0, A[ }\left(\left\|\beta^{2}(t, a, x)\right\|+\|\nabla \beta\|^{2}(t, a, x) d a\right) ; \\ \exists \delta \in(0, A) \text { s.t. } \beta(a, .,)=0 \text { for a } \in(\delta, A) ;\end{array}\right.$
$(H 2) \quad: \quad \mu \in C([0, T] \times[0, A] \times \bar{\Omega}), \mu(t, a, x) \geq 0$ a.e in $\quad Q$

$$
(H 3):\left\{\begin{array}{rc}
\forall t, 0<t<A, & \forall x \in \Omega, \lim _{a \longrightarrow A}^{A} \mu(\iota, a-t+\iota, x) d \iota=+\infty \\
\forall t, A<t<T, & \forall x \in \Omega, \lim _{a \rightarrow 0}^{a} \mu(t-a+\alpha, \alpha, x) d \alpha=+\infty \\
\nabla \mu \in\left[L^{\infty}(Q)\right]^{n}
\end{array}\right.
$$

We also assume that:

- $y^{0}$ and $\hat{y}^{0}$ belong to $L^{2}\left(Q_{A}\right), \xi$ and $\hat{\xi}_{i}$ belong to $L^{2}(\Sigma)$,
- the reals $\tau, \lambda_{i} 1 \leq i \leq M$ are sufficiently small and $\left\|\hat{y}^{0}\right\|_{L^{2}\left(Q_{A}\right)} \leq 1$, and we set $\lambda=\left(\lambda_{1}, \ldots, \lambda_{M}\right)$.

Under the above assumptions on the data, one can prove as in [17] that problem (1) has a unique solution in $L^{2}(Q)$. For the sake of simplicity, we denote

$$
\begin{equation*}
y(t, a, x ; \lambda, \tau) \tag{2}
\end{equation*}
$$

the unique solution of (1). Therefore, the map

$$
\begin{equation*}
(\lambda, \tau) \mapsto y(\lambda, \tau) \text { is in } C^{1}\left(\mathbb{R} \times \mathbb{R} ; L^{2}(Q)\right) . \tag{3}
\end{equation*}
$$

For more literature on the model describing the dynamics of population with age dependence and spatial structure as well as for some existence results on such problem, we refer for instance to $[1,3,8,17]$ and the reference therein. Recently S. Sawadogo [16] use the sentinel method to control the migration of a single species population subjected to a migratory phenomenon.

For the model (1), we are interested in identifying the parameters $\lambda_{i}$ without any attempt at computing $\tau \hat{y}^{0}$.

To identify these parameters, we use the theory of sentinel in a general framework. More precisely, Let $O$ be a nonempty open subset of $\Gamma \backslash \Gamma_{1}$ and let $y=y(t, a, x ; \lambda, \tau)=$ $y(\lambda, \tau)$ be the solution of (1). Then for any non-empty open subset $\gamma$ of $\Gamma \backslash \Gamma_{1}$ such that $O \cap \gamma \neq \emptyset$, we look for a function $S(\lambda, \tau)$ solution to the following problem : given $h_{0} \in L^{2}(U \times O)$, find $w \in L^{2}(U \times \gamma)$ such that
i) the function $S$ defined by

$$
\begin{equation*}
S(\lambda, \tau)=\int_{U} \int_{O} h_{0} \frac{\partial y}{\partial \nu}(\lambda, \tau) d t d a d \Gamma+\int_{U} \int_{\gamma} w \frac{\partial y}{\partial \nu}(\lambda, \tau) d t d a d \Gamma, \tag{4}
\end{equation*}
$$

satisfies :

- $S$ is stationary to the first order with respect to missing term $\tau \hat{y}^{0}$

$$
\begin{equation*}
\frac{\partial S}{\partial \tau}(0,0)=0 \quad \forall \hat{y}^{0} \tag{5}
\end{equation*}
$$

- $S$ is sensitive to the first order with respect to pollution terms $\lambda_{i} \hat{\xi}_{i}$ :

$$
\begin{equation*}
\frac{\partial S}{\partial \lambda_{i}}(0,0)=c_{i} \quad 1 \leq i \leq M \tag{6}
\end{equation*}
$$

where $c_{i}, 1 \leq i \leq M$, are given constants not all identically zero.
ii) The control $w$ is of minimal norm in $L^{2}(U \times \gamma)$ among " the admissible controls", i.e.

$$
\begin{equation*}
\|w\|_{L^{2}(U \times \gamma)}^{2}=\min _{\bar{w} \in E}\|\tilde{w}\|_{L^{2}(U \times \gamma)}^{2} \tag{7}
\end{equation*}
$$

where

$$
\begin{equation*}
E=\left\{\tilde{w} \in L^{2}(U \times \gamma), \text { such that }(\tilde{w}, S(\tilde{w})) \text { satisfies }(4)-(7)\right\} \tag{8}
\end{equation*}
$$

Remark 1. J.L.Lions refers to the function $S$ as a sentinel with given sensitivity $c_{i}$. In (6), the $c_{i}$ are chosen according to the importance which is conferred to the component $\xi_{i}$ of the pollution.

Remark 2. Notice that for the J.L.Lions's sentinels defined by (4)-(7), the observatory $O \subset\left(\Gamma \backslash \Gamma_{1}\right)$ is also the support of the control function $w$.

For more information on the theory of sentinel, we refer to [9-11, 14, 15, 20] and the reference therein. We set $y_{0}=y(0,0) \in L^{2}(Q)$, the solution of $(1)$ when $\lambda=0$ and $\tau=0$ and we denote respectively by $y_{\tau}$ and $y_{\lambda_{i}}$, the derivatives of $y$ at $(0,0)$ with respect to $\tau$ and $\lambda_{i}$, i.e. :

$$
y_{\tau}=\lim _{\tau \rightarrow 0} \frac{y(0, \tau)-y(0,0)}{\tau}
$$

and

$$
y_{\lambda_{i}}=\lim _{\lambda_{i} \rightarrow 0} \frac{y\left(\lambda_{i}, 0\right)-y(0,0)}{\lambda}
$$

Then $y_{\tau}$ and $y_{\lambda_{i}}$ are respectively solutions of

$$
\left\{\begin{array}{cccc}
\frac{\partial y_{\tau}}{\partial t}+\frac{\partial y_{\tau}}{\partial a}-\Delta y_{\tau}+\mu y_{\tau} & = & 0 & \text { in } Q  \tag{9}\\
y_{\tau}(0, a, x) & = & \hat{y}^{0} & \text { in } Q_{A} \\
y_{\tau}(t, 0, x) & =\int_{0}^{A} \beta(t, a, x) y_{\tau}(t, a, x) d a & \text { in } Q_{T} \\
y_{\tau} & = & 0 & \text { on } \Sigma
\end{array}\right.
$$

and

$$
\left\{\begin{array}{clcl}
\frac{\partial y_{\lambda_{i}}}{\partial t}+\frac{\partial y_{\lambda_{i}}}{\partial a}-\Delta y_{\lambda_{i}}+\mu y_{\lambda_{i}} & = & 0 & \text { in } \quad Q,  \tag{10}\\
y_{\lambda_{i}}(0, a, x) & = & 0 & \text { in } Q_{A}, \\
y_{\lambda_{i}}(t, 0, x) & =\int_{0}^{A} \beta y_{\lambda_{i}}(t, a, x) d a & \text { in } Q_{T}, \\
y_{\lambda_{i}} & = & \hat{\xi}_{i} \chi_{\Sigma_{1}} & \text { on } \quad \Sigma,
\end{array}\right.
$$

where $\chi_{X}$ denote now and in the sequel, the characteristic function of the set X . Under the assumptions $\left(H_{1}\right)-\left(H_{4}\right)$, the systems (9) and (10) have respectively a unique solution $y_{\tau} \in$ $L^{2}(Q)$ and $y_{\lambda_{i}} \in L^{2}(Q)$ (see $[8,17]$ ). From now on, we make the following assumptions:

- The functions

$$
\begin{equation*}
\hat{\xi}_{i \cdot \chi_{\Sigma_{1}}}, 1 \leq i \leq M \text { are linearly independent } \tag{11}
\end{equation*}
$$

- Any function $\rho$ such that

$$
\left\{\begin{array}{cccc}
\frac{\partial \rho}{\partial t}+\frac{\partial \rho}{\partial a}-\Delta \rho+\mu \rho & = & 0 & \text { in } \tag{12}
\end{array},\right.
$$

is identically zero.
and we set

$$
\begin{equation*}
Y=\operatorname{Span}\left\{\frac{\partial y_{\lambda_{1}}}{\partial \nu} \chi_{\gamma}, \ldots, \frac{\partial y_{\lambda_{M}}}{\partial \nu} \chi_{\gamma}\right\} \tag{13}
\end{equation*}
$$

The vector subspace of $L^{2}(U \times \gamma)$ generated by $M$ functions $\left\{\frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}\right\}_{i=1}^{M}$.

$$
Y_{\theta}=\frac{1}{\theta} Y
$$

The vector subspace of $L^{2}(U \times \gamma)$ generated by $M$ functions $\left\{\frac{1}{\theta} \frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}\right\}_{i=1}^{M}$, where $\theta$ is the positive function precisely defined later on by (31).

Remark 3. We will prove in Lemma 1 that the function $\left\{\frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}\right\}_{i=1}^{M}$ and $\left\{\frac{1}{\theta} \frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}\right\}_{i=1}^{M}$ are linearly independent.

We now consider the following boundary null-controllability problem :
given $h_{0} \in L^{2}(U \times O), w_{0} \in Y_{\theta}$, find $v \in L^{2}(U \times \gamma)$ such that

$$
\begin{equation*}
v \in Y^{\perp} \tag{14}
\end{equation*}
$$

and if $q=q(t, a, x ; v)$ is solution of

$$
\left\{\begin{array}{ccccc}
-\frac{\partial q}{\partial t}-\frac{\partial q}{\partial a}-\Delta q+\mu q & = & \beta q(t, 0, x) & \text { in } & Q,  \tag{15}\\
q & = & h_{0} \chi_{O}+\left(w_{0}-v\right) \chi_{\gamma} & \text { on } & \Sigma, \\
q(T, a, x) & = & 0 & \text { in } & Q_{A}, \\
q(t, A, x) & = & 0 & \text { in } & Q_{T},
\end{array}\right.
$$

$q$ satisfy

$$
\begin{equation*}
q(0, a, x ; v)=0 \quad \text { in } \quad Q_{A} . \tag{16}
\end{equation*}
$$

Remark 4. Let us notice that if $v$ exists, the set

$$
\begin{equation*}
\mathcal{E}=\left\{\bar{v} \in Y^{\perp} \text { such } \quad \text { that } \quad(\bar{v}, \bar{q}=q(t, a, x ; \bar{v})) \text { satisfies } \quad(15)-(16)\right\} \tag{17}
\end{equation*}
$$

is a non-empty closed, and convex set in $L^{2}(U \times \gamma)$. Therefore there exists $v \in \mathcal{E}$ of minimal norm.

The problem (14) - (16) is a null boundary controllability problem with constraint on the control. When $Y^{\perp}=L^{2}(U \times \gamma)$, this problem becomes a null controllability problem without constraint on the control. This kind of problem has been studied by many authors with various methods [2,5]. In this paper we solve the boundary null controllability problem with constraint on the control (14) - (16), this allows us to prove the existence of the sentinel with given sensitivity (4) - (7). More precisely, we have the following results:

Theorem 1. Let $\Omega$ be a bounded open subset of $\mathbb{R}^{N}$ with boundary $\Gamma$ of class $C^{\infty}$. Let $\Gamma_{1}$ be a non-empty open subset of $\Gamma$. Let also $O$ and $\gamma$ be two non empty subsets of $\Gamma \backslash \Gamma_{1}$, such that $O \cap \gamma \neq 0$. Assume that the assumptions of the data of the system (1) are satisfied. Assume also that (11) and (12) holds. Then the existence of sentinel (4) - (7) holds if and only if, the boundary null-controllability problem with constraints on the control (14) - (16) has a solution.

To prove the boundary null-controllability problem with constraints on the control (14) - (16), we use an inequality of Carleman adapted to the constraint that we establish by means of a global Carleman inequality. More precisely we prove the following results.

Theorem 2. Assume that the hypotheses of Theorem 1 are satisfied. Then there exists a positive real weight function $\theta$ (a precise definition of $\theta$ will be given later on (31)) such that, for any function $h_{0} \in L^{2}(U \times O)$ with $\theta h_{0} \in L^{2}(U \times O)$ there exists a unique control $\hat{v} \in L^{2}(U \times \gamma)$ such that $(\hat{v}, \hat{q})$ with $\hat{q}=q(\hat{v})$ is solution of null boundary controllability problem with constraint on the control (14) - (16) and provides a control $\hat{w}=w_{0} \chi_{\gamma}-\hat{v}$ of the sentinel problem satisfying (7) . Moreover, the control $\hat{w}$ is given by

$$
\begin{equation*}
\hat{w}=P\left(w_{0}\right)+(I-P)\left(\frac{\partial \hat{\rho}}{\partial \nu} \chi_{\gamma}\right), \tag{18}
\end{equation*}
$$

where $P$ is the orthogonal projection operator from $L^{2}(U \times \gamma)$ into $Y, w_{0} \in Y_{\theta}$ depends on $h_{0}$ and $c_{i}, i \in\{1, \ldots, M\}$, and will be precisely determined in (27) and $\hat{\rho}$ satisfies

$$
\left\{\begin{array}{ccccc}
\frac{\partial \hat{\rho}}{\partial t}+\frac{\partial \hat{\rho}}{\partial a}-\Delta \hat{\rho}+\mu \hat{\rho} & = & 0 & \text { in } & Q  \tag{19}\\
\hat{\rho} & = & 0 & \text { on } & \Sigma \\
\hat{\rho}(t, 0, x) & = & \int_{0}^{A} \beta(t, a, x) \hat{\rho}(t, a, x) d a & \text { in } \quad Q_{T}, \\
\hat{\rho}(0, ., .) & = & 0 & \text { in } \quad Q_{A},
\end{array}\right.
$$

The rest of the paper is organized as follows: section 2 is devoted to the equivalence between the sentinel problem and the null boundary controllability problem with constraint on the control. In this section we give the proof of Theorem 1. In section 3, we establish Carleman inequalities necessary to solve the boundary null-controllability problem with constraint on the control (14) - (16). In subsection 3.2 we give the proof of Theorem 2. In section 4, we formulate the sentinel and we identify the parameters.

## 2. Equivalence between the sentinel problem and the null boundary controllability problem with constraint on the control

In this subsection we prove Theorem 1. But before going further, we need the following result:
Lemma 1. Assume that (11) and (12) holds. Then the functions $\frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}, 1 \leq i \leq M$ are linearly independent. Moreover the functions $\frac{1}{\theta} \frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}, 1 \leq i \leq M$, are also linearly independent.

Proof. Let $\alpha_{i} \in \mathbb{R}, 1 \leq i \leq M$ be such that $\sum_{i}^{M} \alpha_{i} \frac{\partial y_{\lambda_{i}}}{\partial \nu} \chi_{\gamma}=0$. Set $k=\sum_{i}^{M} \alpha_{i} y_{\lambda_{i}}$, using (10) $k$ is solution of

$$
\left\{\begin{array}{ccccc}
\frac{\partial k}{\partial t}+\frac{\partial k}{\partial a}-\Delta k+\mu k & = & 0 & \text { in } & Q  \tag{20}\\
k(0, a, x) & = & 0 & \text { in } & Q_{A}, \\
k(t, 0, x) & = & \int_{0}^{A} \beta k(t, a, x) d a & \text { in } & Q_{T}, \\
k & = & \sum_{i=1}^{M} \alpha_{i} \hat{\xi}_{i} \cdot \chi_{\Sigma_{1}} & \text { on } & \Sigma, \\
\frac{\partial k}{\partial \nu} & = & 0 & \text { on } & U \times \gamma .
\end{array}\right.
$$

Assumption (12) allows us to say that $k=0$ in $Q$. Therefore, we deduce that $\sum_{i=1}^{M} \alpha_{i} \widehat{\xi}_{i} \cdot \chi_{\Sigma_{1}}=$ 0 on $\Sigma$. Then it follows from (11) that $\alpha_{i}=0$ for $1 \leq i \leq M$. The second assertion of the lemma follows immediately.

Now, let us prove Theorem 1. To this end, we interpret (5) and (6). Actually, in view of (4), the stationary condition (5) and respectively the sensitivity conditions (6) hold if and only if

$$
\begin{equation*}
\int_{U} \int_{O} h_{0} \frac{\partial y_{\tau}}{\partial \nu} d t d a d \Gamma+\int_{U} \int_{\gamma} w \frac{\partial y_{\tau}}{\partial \nu} d t d a d \Gamma=0 . \forall \hat{y}^{0},\left\|\hat{y}^{0}\right\|_{L^{2}\left(Q_{A}\right)} \leq 1 \tag{21}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{U} \int_{O} h_{0} \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma+\int_{U} \int_{\gamma} w \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma=c_{i}, 1 \leq i \leq M \tag{22}
\end{equation*}
$$

Therefore, in order to transform equation (21), we consider the following adjoint equation

$$
\left\{\begin{array}{ccccc}
-\frac{\partial q}{\partial t}-\frac{\partial q}{\partial a}-\Delta q+\mu q & = & \beta q(t, 0, x) & \text { in } & Q  \tag{23}\\
q & = & h_{0} \chi_{O}+w \chi_{\gamma} & \text { on } & \Sigma \\
q(T, a, x) & = & 0 & \text { in } & Q_{A} \\
q(t, A, x) & = & 0 & \text { in } & Q_{T}
\end{array}\right.
$$

Since $h_{0} \chi_{O}+w \chi_{\gamma} \in L^{2}(\Sigma)$, the assumptions $(H 1)-(H 2)$ ensure that that (23) has a unique solution $q \in L^{2}(Q)$. Now multiplying both sides of the differential equation in (23) by $y_{\tau}$ solution of (9) and integrating by parts in $Q$, we get

$$
\begin{equation*}
\int_{U} \int_{O} h_{0} \frac{\partial y_{\tau}}{\partial \nu} d t d a d \Gamma+\int_{U} \int_{\gamma} w \frac{\partial y_{\tau}}{\partial \nu} d t d a d \Gamma=\int_{0}^{A} \int_{\gamma} q(0, a, x) \hat{y}^{0} d a d x \forall \hat{y}^{0} \in L^{2}\left(Q_{A}\right) \tag{24}
\end{equation*}
$$

Thus, the condition (5) or (21) holds if and only if

$$
\begin{equation*}
q(0, a, x ; v)=0 \quad \text { in } Q_{A} \tag{25}
\end{equation*}
$$

Then, multiplying both sides of the differential equation in (23) by $y_{\lambda_{i}}$ solution of (10) and integrating by parts in $Q$, we have

$$
\int_{U} \int_{O} h_{0} \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma+\int_{U} \int_{\gamma} w \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma=\int_{\Sigma_{1}} \frac{\partial q^{\partial \nu}}{\partial \nu} \hat{\xi}_{i} \cdot \chi_{\Gamma_{1}} d t d a, \quad 1 \leq i \leq M
$$

Thus, the condition the condition (6) or (22) is equivalent to

$$
\begin{equation*}
\int_{\Sigma_{1}} \frac{\partial q}{\partial \nu} \hat{\xi}_{i} \cdot \chi_{\Gamma_{1}} d t d a=c_{i}, \quad 1 \leq i \leq M \tag{26}
\end{equation*}
$$

Now, consider the matrix

$$
\left(\int_{0}^{T} \int_{0}^{A} \int_{\gamma} \frac{1}{\theta} \frac{\partial y_{\lambda_{j}}}{\partial \nu} \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma\right)_{1 \leq i, j \leq M}
$$

Since this matrix is symmetric positive definite therefore, there exists a unique $w_{0} \in Y_{\theta}$ such that

$$
\begin{equation*}
c_{i}-\int_{U} \int_{O} h_{0} \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma=\int_{U} \int_{\gamma} w_{0} \frac{\partial y_{\lambda_{i}}}{\partial \nu} d t d a d \Gamma . \quad 1 \leq i \leq M \tag{27}
\end{equation*}
$$

Consequently, combining (27) with (22), we observe that condition (6) ( or the constraints (26)) holds if and only if

$$
w-w_{0}=-v \in Y^{\perp}
$$

where Y is given by (13). Replacing $w$ by $w_{0}-v$ in the second expression of (23), we obtain (15). We just have proved that the sentinel problem (4) - (7) hold if and only if null controllability problem with constraint on the control (14) - (16) has a solution.

Remark 5. If $\mathcal{E}$ is the set of admissible control $v \in L^{2}(U \times \gamma)$ such that (14)-(16) is satisfied, then $\mathcal{E}$ is a closed convex subset of $L^{2}(U \times \gamma)$. Since $w_{0}-\mathcal{E}$ is also a closed convex subset of $L^{2}(U \times \gamma)$, we can obtain $w$ to be of minimum norm in $L^{2}(U \times \gamma)$ by minimizing the norm of $w_{0}-v$ when $v \in \mathcal{E}$. Then the pair $(v, q(v))$ satisfying (14) - (16) necessarily provides a control $w$ satisfying (7)

## 3. Study of the boundary null-controllability problem with constraint on the control

In this section, we prove existence of the solution of the boundary null controllability problem (14) - (16) and of course uniqueness if we want the control to be of minimal norm among admissible controls. The main tool we use is an observability inequality adapted to the constraint (14) which itself is a consequence of a global Carleman inequality.

### 3.1. An adapted Carleman inequality

The observability inequality we are looking for is a consequence of the global Carleman's inequality. We consider an auxiliary function an auxillary function $\psi \in C^{2}(\bar{\Omega})$ which satisfies the following conditions :

$$
\begin{align*}
\psi(x) & >0, \forall x \in \Omega \\
\nabla \psi & >\alpha, \forall x \in \bar{\Omega}, \\
\psi(x) & =0, \forall x \in \Gamma \backslash \gamma,  \tag{28}\\
\frac{\partial \psi}{\partial \nu} & <0, \forall x \in \Gamma \backslash \gamma .
\end{align*}
$$

Such a function exists according to A. Fursikov and O. Yu. Imanuvilov [7]. For any positive parameter value $\lambda$ we define the following weight functions :

$$
\begin{equation*}
\varphi(t, a, x)=\frac{e^{\lambda\left(m|\psi|_{\infty}+\psi(x)\right)}}{a t(A-a)(T-t)}, \tag{29}
\end{equation*}
$$

$$
\begin{equation*}
\eta(t, a, x)=\frac{e^{2 \lambda m|\psi|_{\infty}}-e^{\lambda\left(m|\psi|_{\infty}+\psi(x)\right)}}{a t(A-a)(T-t)} \tag{30}
\end{equation*}
$$

with $m \geq 1$. Since $\varphi$ does not vanish in Q , for all $s>0$ and $\lambda>0$, we set

$$
\begin{equation*}
\frac{1}{\theta^{2}}=\min \left[e^{-2 s \eta}\left(\varphi^{-1}, \varphi, \varphi^{3}, \varphi,\left|\frac{\partial \psi}{\partial \nu}\right|\right)\right] \tag{31}
\end{equation*}
$$

and we adopt the following notations :

$$
\left\{\begin{align*}
L & =\frac{\partial}{\partial t}+\frac{\partial}{\partial a}-\Delta+\mu I  \tag{32}\\
L^{*} & =-\frac{\partial}{\partial t}-\frac{\partial}{\partial a}-\Delta+\mu I \\
\mathcal{V} & =\left\{\rho \in C^{\infty}(\bar{Q}), \rho=0 \text { on } \Sigma\right\}
\end{align*}\right.
$$

Using the notations given by (32) and the definition of $\theta$ given by (31), we have the following boundary Carleman inequality:

Proposition 1. [Global Carleman inequality] Let $\psi, \varphi$ and $\eta$ be defined respectively by $(28)-(30)$. Then, there exists numbers $\lambda_{0}=\lambda_{0}(\gamma, \mu)>1, s_{0}=s_{0}(\gamma, \mu, T)>1, C_{0}=$ $C_{0}(\gamma, \mu)>0$ and $C_{1}=C_{1}(\gamma, \mu)>0$ such that for any $\lambda \geq \lambda_{0}$, for any $s \geq s_{0}$, for any $\rho \in \nu$, the following estimate holds :

$$
\begin{gather*}
\int_{Q} \frac{e^{-2 s \eta}}{s \varphi}\left(\left|\rho_{t}+\rho_{a}\right|^{2}+|\Delta \rho|^{2}\right) d a d t d x+\int_{Q} e^{-2 s \eta}\left(s \lambda^{2} \varphi|\nabla \rho|^{2}+s^{3} \lambda^{4} \varphi^{3}|\rho|^{2}\right) d t d a d x \\
+C_{0} \int_{0}^{T} \int_{0}^{A} \int_{\Gamma \backslash \gamma} s e^{-2 s \eta} \varphi\left(-\frac{\partial \psi}{\partial \nu}\right)\left|\frac{\partial \rho}{\partial \nu}\right|^{2} d t d a d \Gamma \\
\leq C_{1}\left[\int_{Q} e^{-2 \eta}|L \rho|^{2} d t d a d x+\int_{0}^{T} \int_{0}^{A} \int_{\gamma} s e^{-2 \eta} \varphi\left|\frac{\partial \rho}{\partial \nu}\right|^{2} d t d a d \Gamma\right] \tag{33}
\end{gather*}
$$

Proof. See [18]
As $\psi$ belong to $C^{2}(\bar{\Omega})$ and $\varphi e^{-2 s \eta}$ is bounded, then $\frac{1}{\theta}$ is also bounded in $Q$. Hence, from Proposition 1, we have this other inequality :

Proposition 2. Let $\theta$ be defined by (31). Then, there exists numbers $\lambda_{0}=\lambda_{0}(\Omega, \gamma, \mu)>1$, $s_{0}=s_{0}(\Omega, \gamma, \mu, T)>1, C_{0}=C_{0}(\Omega, \gamma, \mu)>0$, and $C_{1}=C_{1}(\Omega, \gamma, \mu)>0$ such that, for any $\lambda \geq \lambda_{0}$, for any $s \geq s_{0}$, and for any $\rho \in \mathcal{V}$,

$$
\begin{align*}
\int_{U} \int_{\Omega} \frac{1}{\theta^{2}}\left(\left|\frac{\partial \rho}{\partial \nu}\right|^{2}+|\Delta \rho|^{2}\right. & \left.+|\nabla \rho|^{2}+|\rho|^{2}\right) d t d a d \Gamma+C_{0} \int_{U} \int_{\Gamma} \frac{1}{\theta^{2}}\left|\frac{\partial \rho}{\partial \nu}\right|^{2} d t d a d \Gamma \\
& \leq C_{1}\left[\int_{Q}|L \rho|^{2} d t d a d x+\int_{U} \int_{\gamma}\left|\frac{\partial \rho}{\partial \nu}\right|^{2} d t d a d \Gamma\right] \tag{34}
\end{align*}
$$

Lemma 2. Under the assumptions of Lemma 1. Let $Y$ be the real vector subspace of $L^{2}(U \times \gamma)$ of finite dimension defined in (13). Then any function $\rho$ such that

$$
\left\{\begin{array}{ccccc}
\frac{\partial \rho}{\partial t}+\frac{\partial \rho}{\partial a}-\Delta \rho+\mu \rho & =0 & \text { in } & Q  \tag{35}\\
\rho(0, ., .) & =0 & \text { in } & Q_{A}, \\
\rho & = & 0 & \text { on } & \Sigma \backslash \Sigma_{1}, \\
\left.\frac{\partial \rho}{\partial \nu}\right|_{\gamma} \in Y, & & &
\end{array}\right.
$$

is identically zero.
Proof. For any $\rho$ verifying (35) there exists $\alpha_{i} \in \mathbb{R}, 1 \leq i \leq M$, such that $\frac{\partial \rho}{\partial \nu}=$ $\sum_{i=1}^{M} \alpha_{i} \frac{\partial y_{i}}{\partial \nu}$. We set $z=\rho-\sum_{i=1}^{M} \alpha_{i} y_{i}$. Using (10), we have

$$
\left\{\begin{array}{ccccc}
\frac{\partial z}{\partial t}+\frac{\partial z}{\partial a}-\Delta z+\mu z & = & \text { in } & Q,  \tag{36}\\
z(0, ., .) & = & 0 & \text { in } & Q_{A}, \\
z & = & 0 & \text { on } & \Sigma \backslash \Sigma_{1}, \\
\frac{\partial z}{\partial \nu} & = & 0 & \text { on } & U \times \gamma .
\end{array}\right.
$$

As $\gamma \subset \Gamma \backslash \Gamma_{1}$, we have $z=0$ and $\frac{\partial z}{\partial \nu}=0$ in $U \times \gamma$. Then it follows from (12) that $z=0$ in $Q$. Consequently, we deduce on the one hand that $\rho=\sum_{i=1}^{M} \alpha_{i} y_{i}$ and on the other hand that $\sum_{i=1}^{M} \alpha_{i} \widehat{\xi}_{i}=0$ on $\Sigma_{1}$. Hence, it follows from assumption (11) that $\alpha_{i}=0$ for $1 \leq i \leq M$.Thus, $\rho=0$ in $Q$.

Proposition 3 (Adapted Carleman inequality). Under the Assumption of Lemma1 . Let $Y$ be the real vector subspace of $L^{2}(U \times \gamma)$ of finite dimension defined in (13) and $P$ be the orthogonal projection operator from $L^{2}(U \times \gamma)$ into $Y$. Let also $\theta$ be the function defined by (31). Then, there exists numbers $\lambda_{0}=\lambda_{0}(\Omega, \gamma, \mu)>1, s_{0}=s_{0}(\Omega, \gamma, \mu, T)>1$, $C_{0}=C_{0}(\Omega, \gamma, \mu)>0$ and $C_{1}=C_{1}(\Omega, \gamma, \mu)>0$ such that, for any $\lambda \geq \lambda_{0}$, for any $s \geq s_{0}$, and for any $\rho \in \mathcal{V}$,

$$
\begin{equation*}
\int_{U} \int_{\Omega} \frac{1}{\theta^{2}}\left|\frac{\partial \rho}{\partial \nu}\right|^{2} d t d a d \Gamma \leq C_{1}\left[\int_{Q}|L \rho|^{2} d t d a d x+\int_{U} \int_{\gamma}\left|P \frac{\partial \rho}{\partial \nu}-\frac{\partial \rho}{\partial \nu}\right|^{2} d t d a d \Gamma\right] . \tag{37}
\end{equation*}
$$

Proof. As in [9], we use a well known compactness-uniqueness argument and the inequality (34). Indeed, suppose that (37) does not hold. Then for any $j \in \mathbb{N}$, there exists $\rho_{j} \in \mathcal{V}$ such that

$$
\begin{equation*}
\int_{U} \int_{\Omega}\left|L \rho_{j}\right|^{2} d t d a d x \leq \frac{1}{j}, \tag{38}
\end{equation*}
$$

$$
\begin{gather*}
\int_{U} \int_{\gamma}\left|P \frac{\partial \rho_{j}}{\partial \nu}-\frac{\partial \rho_{j}}{\partial \nu} \chi_{\gamma}\right|^{2} d t d a d \Gamma \leq \frac{1}{j}  \tag{39}\\
\int_{U} \int_{\Gamma} \frac{1}{\theta^{2}}\left|\frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma=1 \tag{40}
\end{gather*}
$$

In what follows, we prove in three steps that (38) - (40) yields contradiction.
Step 1. We have

$$
\begin{align*}
\int_{U} \int_{\gamma} \frac{1}{\theta^{2}}\left|P \frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma & \leq 2 \int_{U} \int_{\gamma} \frac{1}{\theta^{2}}\left|P \frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma \\
& +2 \int_{U} \int_{\gamma} \frac{1}{\theta^{2}}\left|P \frac{\partial \rho_{j}}{\partial \nu}-\frac{\partial \rho_{j}}{\partial \nu} \chi_{\gamma}\right|^{2} d t d a d \Gamma \tag{41}
\end{align*}
$$

Since $\frac{1}{\theta^{2}}$ is bounded, using (38) and (39), it follows that there exists a positive constant $C$ such that

$$
\begin{equation*}
\forall j \in \mathbb{N}, \quad \int_{U} \int_{\gamma}\left|P \frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma \leq C \tag{42}
\end{equation*}
$$

As $\frac{\partial \rho}{\partial \nu} \chi_{\gamma}=P \frac{\partial \rho}{\partial \nu} \chi_{\gamma}+\left(\frac{\partial \rho}{\partial \nu} \chi_{\gamma}-P \frac{\partial \rho}{\partial \nu} \chi_{\gamma}\right)$, using (40) and (42), we obtain

$$
\begin{equation*}
\left\|\frac{\partial \rho_{j}}{\partial \nu}\right\|_{L^{2}(U \times \gamma)}^{2} \leq C \tag{43}
\end{equation*}
$$

Step 2. Let $L^{2}\left(\frac{1}{\theta}, U \times \gamma\right)=\left\{\rho \in L^{2}(U \times \Omega) ; \int_{U} \int_{\Gamma} \frac{1}{\theta^{2}}\left|\frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma<\infty\right\}$.
Then in view of (40) and (43), we deduce from (34) that, $\left(\frac{\partial \rho}{\partial t}+\frac{\partial \rho}{\partial a}\right),\left(\frac{\partial \rho_{j}}{\partial \nu}\right),\left(\nabla \rho_{j}\right),\left(\rho_{j}\right)$ and $\left(\Delta \rho_{j}\right)$ are bounded in $L^{2}\left(\frac{1}{\theta}, U \times \gamma\right)$. Let us the take a subsequence still denoted by $\left(\rho_{j}\right)$ such that

$$
\begin{align*}
\rho_{j} & \rightharpoonup \rho \quad \text { weakly in } L^{2}\left(\frac{1}{\theta}, U \times \gamma\right)  \tag{44}\\
\frac{\partial \rho_{j}}{\partial \nu} & \rightharpoonup \frac{\partial \rho}{\partial \nu} \quad \text { weakly in } L^{2}\left(\frac{1}{\theta}, U \times \gamma\right) \tag{45}
\end{align*}
$$

Then follows from $(28)-(30)$ and the definition of $\frac{1}{\theta}$ given by $(31)$ that $\left(\rho_{j}\right)$ and $\left(\Delta \rho_{j}\right)$ are bounded in $L^{2}(] \beta, T-\beta[\times] \alpha, A-\alpha[\times \Omega)$ for any $\beta>0$ and any $\alpha>0$. In particular, for all $\beta>0$ and any $\alpha>0$, we have

$$
\begin{gathered}
\rho_{j} \rightharpoonup \rho \quad \text { weakly in } L^{2}(] \beta, T-\beta[\times] \alpha, A-\alpha[\times \Omega) \\
\frac{\partial \rho_{j}}{\partial \nu} \rightharpoonup \frac{\partial \rho}{\partial \nu} \quad \text { weakly in } L^{2}(] \beta, T-\beta[\times] \alpha, A-\alpha[\times \Sigma)
\end{gathered}
$$

which implies that

$$
\begin{gathered}
\rho_{j} \rightharpoonup \rho \quad \text { weakly in } D^{\prime}(Q) \\
\frac{\partial \rho_{j}}{\partial \nu} \rightharpoonup \frac{\partial \rho}{\partial \nu} \quad \text { weakly in } D^{\prime}(\Sigma) .
\end{gathered}
$$

Therefore, we get from (38) and (43) that

$$
\begin{gather*}
L \rho_{j} \longrightarrow L \rho=0 \quad \text { strongly in } L^{2}(U \times \Omega),  \tag{46}\\
\frac{\partial \rho_{j}}{\partial \nu} \rightharpoonup \frac{\partial \rho}{\partial \nu} \text { strongly in } L^{2}(U \times \gamma) . \tag{47}
\end{gather*}
$$

And, since $P$ is a compact operator, we deduce from (47) that

$$
\begin{equation*}
P \frac{\partial \rho_{j}}{\partial \nu} \longrightarrow P \frac{\partial \rho}{\partial \nu} \quad \text { strongly in } L^{2}(U \times \gamma) . \tag{48}
\end{equation*}
$$

In view of (39), we also have

$$
\begin{equation*}
\frac{\partial \rho_{j}}{\partial \nu}-P \frac{\partial \rho_{j}}{\partial \nu} \longrightarrow 0 \quad \text { strongly in } L^{2}(U \times \gamma) . \tag{49}
\end{equation*}
$$

Thus combining (48) and (49), we get

$$
\begin{equation*}
P \frac{\partial \rho_{j}}{\partial \nu} \longrightarrow \frac{\partial \rho_{j}}{\partial \nu} \quad \text { strongly in } L^{2}(U \times \gamma) . \tag{50}
\end{equation*}
$$

Thanks to the uniqueness of the limit in $L^{2}(U \times \gamma)$, the convergence relations (48)-(49) and (50) imply that $P \frac{\partial \rho}{\partial \nu}=\frac{\partial \rho}{\partial \nu} \chi_{\gamma}$. This means that $\frac{\partial \rho}{\partial \nu} \chi_{\gamma} \in Y$. We thus have proved that $\rho$ verifies (35). Hence thanks to Lemma $2, \rho$ is identically zero.

Therefore, (50) becomes

$$
\begin{equation*}
\frac{\partial \rho_{j}}{\partial \nu} \longrightarrow 0 \text { strongly in } L^{2}(U \times \gamma) . \tag{51}
\end{equation*}
$$

Step 3. Since $\rho_{j} \in \nu$, it follows from the observability inequality (34) that

$$
\int_{U} \int_{\Omega} \frac{1}{\theta^{2}}\left|\frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma \leq C_{1}\left[\int_{Q}\left|L \rho_{j}\right|^{2} d t d a d x+\int_{U} \int_{\gamma}\left|\frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d \Gamma\right] .
$$

Therefore passing this latter inequality to the limit while using (46) and (51), we obtain

$$
\lim _{j \longrightarrow \infty} \int_{U} \int_{\Omega}\left|\frac{\partial \rho_{j}}{\partial \nu}\right|^{2} d t d a d x=0
$$

The contradiction occurs with (40).

### 3.2. Proof of Theorem 2

In this subsection, we are concerned with the proof of Theorem 2. That is, the optimality system for the control $\hat{v}$ such that the pair $(\hat{v} ; \hat{q})$ verifies (14) - (16). Since a classical way to derive this optimality system is the method of penalization due to J.L.Lions [11], here we use this method.

Step 1. Let $w_{0}$ be defined by (27). If $v \in Y^{\perp}$ and $q$ is solution of (15) then $q(0, \ldots.) \in$ $L^{2}\left(Q_{A}\right)$ and we can define the functional

$$
\begin{equation*}
J_{\epsilon}(v)=\frac{1}{2}\left\|w_{0}-v\right\|_{L^{2}(U \times \gamma)}^{2}+\frac{1}{2 \epsilon}\|q(0, ., .)\|_{L^{2}\left(Q_{A}\right)}^{2} \tag{52}
\end{equation*}
$$

We consider the optimal control problem: Find $v_{\epsilon} \in Y^{\perp}$ such that

$$
\begin{equation*}
J_{\epsilon}\left(v_{\epsilon}\right)=\min _{v \in Y^{\perp}} J_{\epsilon}(v) \tag{53}
\end{equation*}
$$

Since $Y^{\perp}$ is a closed and convex subset of $L^{2}(U \times \gamma)$, it is classical to prove that there exists a unique solution to (53). If we write $q_{\epsilon}$ the solution of (15) corresponding to $v_{\epsilon}$ using an adjoint state $\rho_{\epsilon}$, we have that the triplet
( $q_{\epsilon}, \rho_{\epsilon} v_{\epsilon}$ ) is solution of the first order optimality system:

$$
\begin{align*}
& \left\{\begin{array}{ccccc}
L^{*} q_{\epsilon} & = & \beta q_{\epsilon}(t, 0, x) & \text { in } \quad Q, \\
q_{\epsilon}(T, a, x) & = & 0 & \text { in } Q_{A}, \\
q_{\epsilon}(t, A, x) & = & 0 & \text { in } Q_{T}, \\
q_{\epsilon} & = & h_{0} \chi_{O}+\left(w_{0}-v_{\epsilon}\right) \chi_{\gamma} & \text { on } & \Sigma,
\end{array}\right.  \tag{54}\\
& \left\{\begin{array}{ccccc}
L \rho_{\epsilon} & = & 0 & \text { in } & Q, \\
\rho_{\epsilon}(0, a, x) & = & \frac{1}{\epsilon} q_{\epsilon}(0, a, x) & \text { in } & Q_{A}, \\
\rho_{\epsilon}(t, 0, x) & = & \int_{0}^{A} \beta(t, a, x) \rho_{\epsilon}(t, a, x) d a & \text { in } & Q_{T}, \\
\rho_{\epsilon} & = & 0 & \text { on } & \Sigma,
\end{array}\right.  \tag{55}\\
& v_{\epsilon}=\left(w_{0} \chi_{\gamma}-\frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma}\right)-P\left(w_{0} \chi_{\gamma}-\frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma}\right) \in Y^{\perp} . \tag{56}
\end{align*}
$$

## Step 2.

Multiplying the state equation (54) by $\rho_{\epsilon}$ and integrating by parts over $Q$, we get

$$
\frac{1}{\epsilon}\left\|q_{\epsilon}(0, ., .)\right\|_{L^{2}\left(Q_{A}\right)}^{2}=\int_{U} \int_{O} h_{0} \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma+\int_{U} \int_{\gamma}\left(w_{0}-v_{\epsilon}\right) \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma
$$

Which in view of (56) and the fact that $v_{\epsilon} \in Y^{\perp}$ give

$$
\begin{aligned}
\frac{1}{\epsilon}\left\|q_{\epsilon}(0, ., .)\right\|_{L^{2}\left(Q_{A}\right)}^{2}= & \int_{U} \int_{O} h_{0} \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma \\
& +\int_{U} \int_{\gamma}\left(w_{0}-v_{\epsilon}\right)\left(w_{0}-v_{\epsilon}-P\left(w_{0} \chi_{\gamma}-\frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma}\right)\right) d t d \Gamma \\
= & \int_{U} \int_{O} h_{0} \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma \\
& -\left\|w_{0}-v_{\epsilon}\right\|_{L^{2}(U \times \gamma)}+\left\|P w_{0} \chi_{\gamma}\right\|_{L^{2}(U \times \gamma)}+\int_{U} \int_{\gamma} w_{0} \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma .
\end{aligned}
$$

As on $U \times \gamma$

$$
w_{0}-v_{\epsilon}=P w_{0} \chi_{\gamma}+(I-P) \frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma} .
$$

We have that

$$
\left\|w_{0}-v_{\epsilon}\right\|_{L^{2}(U \times \gamma)}=\left\|(I-P) \frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma}\right\|_{L^{2}(U \times \gamma)}^{2}+\left\|P w_{0} \chi_{\gamma}\right\|_{L^{2}(U \times \gamma)}^{2}
$$

so that

$$
\begin{aligned}
\frac{1}{\epsilon}\left\|q_{\epsilon}(0, ., .)\right\|_{L^{2}\left(Q_{A}\right)}^{2}+\left\|(I-P) \frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma}\right\|_{L^{2}(U \times \gamma)}^{2}= & \int_{U} \int_{O} h_{0} \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma \\
& +\int_{U} \int_{\gamma} w_{0} \frac{\partial \rho_{\epsilon}}{\partial \nu} d t d \Gamma .
\end{aligned}
$$

This implies that

$$
\begin{align*}
\frac{1}{\epsilon}\left\|q_{\epsilon}(0, ., .)\right\|_{L^{2}\left(Q_{A}\right)}^{2}+\left\|(I-P) \frac{\partial \rho_{\epsilon}}{\partial \nu} \chi_{\gamma}\right\|_{L^{2}(U \times \gamma)}^{2} \leq & \left(\int_{U} \int_{O}\left(\theta h_{0}\right)^{2} d t d \Gamma\right)^{\frac{1}{2}}\left(\int_{U} \int_{\gamma} \frac{1}{\theta^{2}} \frac{\partial \rho_{\epsilon}{ }^{2}}{\partial \nu} d t d \Gamma\right)^{\frac{1}{2}} \\
& +\left(\int_{U} \int_{O}\left(\theta w_{0}\right)^{2} d t d \Gamma\right)^{\frac{1}{2}}\left(\int_{U} \int_{\gamma} \frac{1}{\theta^{2}} \frac{\partial \rho_{\epsilon}{ }^{2}}{\partial \nu} d t d \Gamma\right)^{\frac{1}{2}} . \tag{57}
\end{align*}
$$

If we apply the adapted Carleman inequality (37) to $\rho_{\epsilon}$ we obtain

$$
\begin{equation*}
\int_{U} \int_{\Gamma} \frac{1}{\theta^{2}}\left|\frac{\partial \rho_{\epsilon}}{\partial \nu}\right|^{2} d t d \Gamma \leq C \int_{U} \int_{\gamma}\left|(I-P) \frac{\partial \rho_{\epsilon}}{\partial \nu} \chi \gamma\right|^{2} d t d a \Gamma, \tag{58}
\end{equation*}
$$

where $C>0$ is independent of $\epsilon$. From (57), the choice of $w_{0} \in Y_{\theta}$ and the hypothesis on $h_{0}$, we deduce that

$$
\begin{array}{r}
\frac{1}{\epsilon}\left\|q_{\epsilon}(0, ., .)\right\|_{L^{2}\left(Q_{A}\right)}^{2}+\frac{1}{2}\left\|(I-P) \frac{\partial \rho_{\epsilon}}{\partial \nu} \chi \gamma\right\|_{L^{2}(U \times \gamma)}^{2} \\
\quad \leq C\left(\int_{U} \int_{\omega} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{59}
\end{array}
$$

and then

$$
\begin{equation*}
\left\|v_{\epsilon}\right\|_{L^{2}(U \times \omega)}^{2} \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{60}
\end{equation*}
$$

$$
\begin{equation*}
\left\|q_{\epsilon} \chi_{\omega}\right\|_{L^{2}(U \times \gamma)}^{2} \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{61}
\end{equation*}
$$

In view of (58) and (59), we get

$$
\begin{equation*}
\left\|\frac{1}{\theta} \frac{\partial \rho_{\epsilon}}{\partial \nu}\right\|_{L^{2}(\Sigma)} \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{62}
\end{equation*}
$$

and using (59) and the fact that $\frac{1}{\theta}$ is bounded, we have

$$
\left\|\frac{1}{\theta} P \frac{\partial \rho_{\epsilon}}{\partial \nu}\right\|_{L^{2}(U \times \gamma)} \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}}
$$

Therefore, $Y$ being a finite dimensional vector subspace of $L^{2}(U \times \gamma)$, we deduce that

$$
\begin{equation*}
\left\|P \frac{\partial \rho_{\epsilon}}{\partial \nu}\right\|_{L^{2}(U \times \gamma)} \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{63}
\end{equation*}
$$

from which we deduce by using (59) that

$$
\begin{equation*}
\left\|\frac{\partial \rho_{\epsilon}}{\partial \nu}\right\|_{L^{2}(U \times \gamma)} \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{64}
\end{equation*}
$$

Using Proposition 2 , we have that

$$
\begin{align*}
& \int_{U} \int_{\Omega} \frac{1}{\theta^{2}}\left(\left|\frac{\partial \rho_{\epsilon}}{\partial \nu}\right|^{2}+\left|\Delta \rho_{\epsilon}\right|^{2}+\left|\nabla \rho_{\epsilon}\right|^{2}+\left|\rho_{\epsilon}\right|^{2}\right) d t d a d \Gamma \\
& \leq C\left(\int_{U} \int_{\gamma} \theta^{2}\left|w_{0}\right|^{2} d t d a d \Gamma+\int_{U} \int_{O} \theta^{2}\left|h_{0}\right|^{2} d t d a d \Gamma\right)^{\frac{1}{2}} \tag{65}
\end{align*}
$$

## Step 3.

We prove the convergence of $\left(v_{\epsilon}, q_{\epsilon}\right)_{\epsilon}$ and $\rho_{\epsilon}$ towards $\hat{v}, \hat{q}$ and $\rho$ as $\epsilon \longrightarrow 0$. According to (60), (61) and (62) we can extract subsequences of $\left(v_{\epsilon}, q_{\epsilon}\right)_{\epsilon}\left(\operatorname{still}\right.$ called $\left(v_{\epsilon}, q_{\epsilon}\right)_{\epsilon}$ ) such that

$$
\begin{align*}
& v_{\epsilon} \rightharpoonup \widetilde{v} \quad \text { weakly in } L^{2}(U \times \gamma),  \tag{66}\\
& q_{\epsilon} \rightharpoonup \widetilde{q} \quad \text { weakly in } L^{2}\left(U ; H_{0}^{1}(\Omega)\right),  \tag{67}\\
& \frac{1}{\theta} \rho_{\epsilon} \rightharpoonup \widetilde{\rho} \quad \text { weakly in } L^{2}\left(\frac{1}{\theta}, Q\right) . \tag{68}
\end{align*}
$$

As $v_{\epsilon}$ belong to $Y^{\perp}$ which is closed vector subspace of $L^{2}(U \times \gamma)$, we have

$$
\begin{equation*}
\widetilde{v} \in Y^{\perp} \tag{69}
\end{equation*}
$$

The traces $(\widetilde{q}(0, .,),. \widetilde{q}(., 0,)),.(\widetilde{q}(T, \ldots), \widetilde{q}(., A,)$.$) and \frac{\partial \widetilde{q}}{\partial \nu}$ exists and belong respectively to $\left(L^{2}\left(Q_{A}\right)\right)^{2} \times\left(L^{2}\left(Q_{T}\right)\right)^{2}$ and $L^{2}(\Sigma)$ (see [8]).

So, using (66) and (67) while passing (54) to the limit as $\epsilon \longrightarrow 0$, we can prove that
$\widetilde{q}$ is solution of

$$
\left\{\begin{array}{ccccc}
L^{*} \tilde{q} & = & \beta \tilde{q}(t, 0, x) & \text { in } & Q,  \tag{70}\\
\tilde{q}(T, a, x) & = & 0 & \text { in } & Q_{A}, \\
\tilde{q}(t, A, x) & = & 0 & \text { in } & Q_{T}, \\
\tilde{q} & = & h_{0} \chi_{O}+\left(w_{0}-\tilde{v}\right) \chi_{\gamma} & \text { on } & \Sigma,
\end{array}\right.
$$

and it follows from (59) that

$$
\begin{equation*}
q_{\epsilon}(0, \ldots .) \rightharpoonup \widetilde{q}(0, ., .)=0 \text { weakly in } L^{2}(Q) . \tag{71}
\end{equation*}
$$

In view of (69), (70) and (71), $(\widetilde{v}, \widetilde{q})$ verifies the null controllability (14) - (16) and there exists a solution to the boundary null controllability problem. Moreover, it is clear from (68) that $\widetilde{\rho}$ satisfies

$$
\left\{\begin{array}{cccc}
L \widetilde{\rho} & = & 0 & \text { in } \\
\tilde{\rho}(t, 0, x) & =\int_{0}^{A} \beta(t, a, x) \tilde{\rho}(t, a, x) d a & & \text { in } \\
Q_{T}, \\
\widetilde{\rho} & =0 & \text { on } & \Sigma
\end{array}\right.
$$

From (64)

$$
\begin{equation*}
\frac{\partial \rho_{\epsilon}}{\partial \nu} \rightharpoonup \frac{\partial \widetilde{\rho}}{\partial \nu} \quad \text { weakly in } L^{2}(U \times \gamma) \tag{72}
\end{equation*}
$$

We know on the one hand that $(\widetilde{v}, \widetilde{q})$ is solution to null controllability (14) - (16), and on the other other hand that, there exists a unique $\hat{v} \in \varepsilon$ such that $\left(w_{0}-\hat{v}\right)$ is of minimal norm in $L^{2}(U \times \gamma)$. If we denote by $\hat{q}$ the corresponding solution to (15), we have $\hat{q}(0, .,)=$.0 and, as $\widetilde{v} \in \mathcal{E}$,

$$
\frac{1}{2}\left\|w_{0}-v_{\epsilon}\right\|_{L^{2}(U \times \gamma)}^{2} \leq J_{\epsilon}\left(v_{\epsilon}\right) \leq J_{\epsilon}(\hat{v})=\frac{1}{2}\left\|w_{0}-\hat{v}\right\|_{L^{2}(U \times \gamma)}^{2}
$$

and

$$
\frac{1}{2}\left\|w_{0}-\hat{v}\right\|_{L^{2}(U \times \gamma)}^{2} \leq \frac{1}{2}\left\|w_{0}-v_{\epsilon}\right\|_{L^{2}(U \times \gamma)}^{2}
$$

Using (66)

$$
\liminf _{\epsilon \rightarrow 0} \frac{1}{2}\left\|w_{0}-v_{\epsilon}\right\|_{L^{2}(U \times \gamma)}^{2} \geq \frac{1}{2}\left\|w_{0}-\hat{v}\right\|_{L^{2}(U \times \gamma)}^{2}
$$

Hence,

$$
\widetilde{v}=\widehat{v}
$$

and

$$
v_{\epsilon} \rightharpoonup \widetilde{v} \quad \text { strongly in } L^{2}(U \times \gamma) .
$$

Writing $\widetilde{\rho}=\widehat{\rho}$, we obtain

$$
\widehat{v}=(I-P)\left(w_{0} \chi_{\gamma}-\frac{\partial \widehat{\rho}}{\partial \nu} \chi_{\gamma}\right) .
$$

## 4. Formulation of the sentinel with given sensitivity and identification of parameters $\lambda_{i}$

According to Theorem 2, if we replace in (4) $w$ by

$$
\hat{w}=P\left(w_{0}\right)+(I-P)\left(\frac{\partial \hat{\rho}}{\partial \nu} \chi_{\gamma}\right)
$$

the function $S$ defined by

$$
S(\lambda, \tau)=\int_{U} \int_{O} h_{0} \frac{\partial y}{\partial \nu}(\lambda, \tau) d t d a d \Gamma+\int_{U} \int_{\gamma}\left(P\left(w_{0}\right)+(I-P)\left(\frac{\partial \hat{\rho}}{\partial \nu} \chi_{\gamma}\right)\right) \frac{\partial y}{\partial \nu}(\lambda, \tau) d t d a d \Gamma
$$

is such that $(\hat{w}, S(\hat{w}))$ verified the sentinel problem (4)-(7). To estimate the parameters $\lambda_{i}$, one proceeds as follows: assume that the solution of (1) when $\lambda=0$ and $\tau=0$ is known. Then, one has the following information

$$
S(\lambda, \tau)-S(0,0) \approx \sum_{i=1}^{M} \lambda_{i} \frac{\partial S}{\partial \lambda_{i}}(0,0)
$$

Therefore, fixing $i \in\{1, \ldots, M\}$ and choosing

$$
\frac{\partial S}{\partial \lambda_{j}}(0,0)=0 \text { for } j \neq i \text { and } \frac{\partial S}{\partial \lambda_{i}}(0,0)=c_{i}
$$

one obtains the following estimate of the parameter $\lambda_{i}$ :

$$
\lambda_{i} \approx \frac{1}{c_{i}}(S(\lambda, \tau)-S(0,0))
$$

we deduce that

$$
\begin{aligned}
\lambda_{i} \approx & \frac{1}{c_{i}}\left\{\int_{U} \int_{O} h_{0}\left(m_{0}-\frac{\partial y_{0}}{\partial \nu} d t d a d \Gamma\right)\right\} \\
& +\frac{1}{c_{i}}\left\{\int_{U} \int_{\gamma}\left(P\left(w_{0}\right)+(I-P)\left(\frac{\partial \hat{\rho}}{\partial \nu} \chi_{\gamma}\right)\right)\left(m_{0}-\frac{\partial y_{0}}{\partial \nu}\right) d t d a d \Gamma\right\}
\end{aligned}
$$

where $m_{0}$ is a measure of the flux of the population taken on the observatory $O \cup \gamma$ and $y_{0}$ is solution of $(1)$ when $\lambda=0$ and $\tau=0$.

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[^0]:    *Corresponding author.
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    Email addresses: mifiambasoma@yahoo.fr (M. Soma), sawasom@yahoo.fr (S. Sawadogo)

